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### Enhancement of O<sub>2</sub> Dissociation on Au(111) by Adsorbed Oxygen: Implications for Oxidation Catalysis

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Abstract: We show that the dissociation probability of O<sub>2</sub> on the reconstructed, Au(111)-herringbone surface is dramatically increased by the presence of some atomic oxygen on the surface. Specifically, at 400 K the dissociation probability of  $O_2$  on oxygen precovered Au(111) is on the order of  $10^{-3}$ , whereas there is no measurable dissociation on clean Au(111), establishing an upper bound for the dissociation probability of 10<sup>-6</sup>. Atomic oxygen was deposited on the clean reconstructed Au(111)-herringbone surface using electron bombardment of condensed NO<sub>2</sub> at 100 K. The dissociation probability for dioxygen was measured by exposing the surface to <sup>18</sup>O<sub>2</sub>. Temperature programmed desorption (TPD) was used to quantify the amount of oxygen dissociation and to study the stability of the oxygen in all cases. Oxygen desorbs as O<sub>2</sub> in a peak centered at 550 K with pseudo-first-order kinetics; i.e., the desorption peak does not shift with coverage. Our interpretation is that the coverage dependence of the activation energy for dissociation ( $\Delta E_{dis}$ ) and/or preexponential factor ( $v_d$ ) may be responsible for the unusual desorption kinetics, implying a possible energy barrier for O<sub>2</sub> dissociation on Au(111). These results are discussed in the context of Au oxidation chemistry and the relationship to supported Au nanoparticles.

#### Introduction

The behavior of oxygen adsorbed on Au is a topic of considerable current interest because of the discovery that highly dispersed, small gold clusters on metal-oxide supports are active catalysts for low-temperature oxidation of CO and propylene.<sup>1-3</sup> Studies of a range of materials, including Au clusters supported on metal oxides and Au single crystals, have been performed to determine the structure of the Au catalyst and the reaction mechanism relating to their activity.<sup>4–9</sup> Although the structures of oxide supported Au catalysts are well-studied,<sup>4,6</sup> there is still controversy regarding the mechanism for the enhanced reactivity. One barrier to understanding the oxidation chemistry is the ability to study the behavior of oxygen on the Au surface. The major obstacle is adsorption of sufficient oxygen on the surface to study reactivity since most sources of oxygen, e.g. O<sub>2</sub> or NO<sub>2</sub>, do not thermally dissociate with a high enough rate to deposit significant amounts of oxygen on the surface.<sup>10–12</sup> Indeed, the

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dissociation of O<sub>2</sub> to yield oxygen is most likely the key, ratelimiting step in the oxidation of CO and hydrocarbons.

To date, several preparation methods of oxygen preadsorbed Au surface have been addressed. Most recently, Mullins and coworkers have deposited oxygen atoms on Au(111) using a supersonic beam of oxygen atoms produced from a radio frequency generated plasma source.13 Atomic oxygen has also been prepared on Au(111) by exposure of ozone,<sup>14</sup> thermal dissociation of gaseous O<sub>2</sub> using hot filaments,<sup>15</sup> O<sup>+</sup> sputtering,<sup>16</sup> and coadsorption of NO<sub>2</sub> and H<sub>2</sub>O.<sup>17</sup> In this study, we describe a new method for depositing atomic oxygen on Au(111) by electron bombardment of NO<sub>2</sub> condensed on the Au(111) surface.

We find that the oxygen exhibits unusual desorption behavior that may be important in determining the reactivity of oxygen on this surface. Specifically, the oxygen recombines to form gaseous O2 with a peak temperature of 550 K in temperature programmed desorption. The O2 desorption feature exhibits pseudofirst-order kinetics, indicating that it is not due to simple recombination of two oxygen atoms. Given that the production of O<sub>2</sub> is the reverse of dissociation, studies of this process provide further insight into a key step in the oxidation processes on Au.

We also find that the dissociation probability of  $O_2$  is increased by more than 3 orders of magnitude when some atomic oxygen is present on the surface. The initial dissociation

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probability of  $O_2$  is less than  $10^{-6}$  on the clean Au(111)herringbone surface. Once some atomic oxygen is seeded on the surface, the dissociation probability is increased to on the order of  $10^{-3}$  at a surface temperature of 400 K. This remarkable result in combination with imaging studies reported elsewhere<sup>18</sup> indicates that  $O_2$  dissociation is very sensitive to the structure of the surface and, in particular, the creation of small gold islands possibly formed by Au atoms released from the surface. It may also indicate that there are important changes in the electronic structure of the surface presented to an incoming  $O_2$ molecule, as has been suggested by recent theoretical results.<sup>19</sup>

#### **Experimental Section**

All experiments were carried out in a stainless steel ultrahigh vacuum chamber described in detail elsewhere.<sup>20</sup> Briefly, the chamber has a base pressure of  $\sim 6 \times 10^{-10}$  Torr and consists of an X-ray photoelectron spectrometer (PHI ESCA 5300) with a Mg anode X-ray source and a hemispherical electron energy analyzer, a quadrapole mass spectrometer (UTI 100c), and low-energy electron diffraction (LEED) optics.

The reconstructed clean Au(111) surface was prepared by cycles of Ar<sup>+</sup> sputtering (1000 eV, 1.3  $\mu$ A) at 300 K followed by annealing at 900 K for 5 min and 700 K for 1 h. This procedure was repeated until no impurities were detected using X-ray photoelectron spectroscopy (XPS). In addition, the consistent surface structure of Au(111) was confirmed by the presence of satellites in LEED due to the herringbone reconstruction on the Au(111) surface.

Atomic oxygen deposited on Au(111) was prepared by exposure of NO<sub>2</sub> followed by electron bombardment. Specifically, NO<sub>2</sub> (Matheson, anhydrous grade) was exposed directly on Au(111) at 100 K with a pressure rise of  $1 \times 10^{-10}$  Torr for 1 min resulting in multilayer NO<sub>2</sub> on the Au(111) surface. Subsequently, the Au(111) surface with multilayer NO<sub>2</sub> was biased to +100 V and exposed to an electron source (UTI 100c mass spectrometer filament with an emission current of 2 mA) for 1 min. The integrated current flux on the surface was 7  $\mu$ A.

<sup>18</sup>O<sub>2</sub> (Matheson, 99 at. %) was exposed on clean or oxygen precovered Au(111) with a pressure rise of  $1 \times 10^{-7}$  Torr using background dosing for 1 min while maintaining the surface temperature at 400 K. The <sup>18</sup>O<sub>2</sub> flux was estimated via the equation:  $N_{\rm inc} = Pt/$  $(2\pi m k_{\rm B}T)^{1/2}$ , where P, m, and T are pressure, mass, and temperature of  $^{18}\mathrm{O}_2$ , respectively, and *t* is the exposure time. The number of  $^{18}\mathrm{O}_2$  that dissociated ( $N_{ads}$ ) was estimated by measuring the amount of <sup>18</sup>O<sup>16</sup>O produced in TPD relative to the amount of <sup>16</sup>O<sub>2</sub> at saturation coverage (0.4 monolayer) and assuming a Au(111) surface density of  $1.3 \times 10^{15}$ atoms/cm<sup>2</sup>. The <sup>18</sup>O<sub>2</sub> dissociation probability was then calculated from the ratio,  $N_{\rm ads}/N_{\rm inc}$ . It should be noted that the absolute value of dissociation probability may not quantitatively accurate due to the uncalibrated ion gauge; however, the quantitative aspect of our measurement does not affect our final conclusions. Analogous experiments were performed using directed dosing-i.e. with an effusive source directed toward the crystal surface such that the background pressure rise was small. These experiments qualitatively reproduced our result of <sup>18</sup>O<sub>2</sub> dissociation; however, the flux could not be as accurately determined in the directed-dosing geometry. We also performed analogous experiments in a separate vacuum chamber with a base pressure of  $<3 \times 10^{-10}$  Torr and using a separate Au(111) crystal to rule out artifacts in the data shown.

All TPD data were taken by UTI 100c mass spectrometer with an average heating rate of 10 K/s, as described in detail previously.<sup>21</sup> All TPD experiments were performed with the crystal biased at -100 V to avoid electron-induced reaction from the mass spectrometer filament.



*Figure 1.* Temperature programmed desorption of NO<sub>2</sub> and O<sub>2</sub> from the Au(111): after (solid line) and before (dashed line) electron bombardment of condensed NO<sub>2</sub>. No O<sub>2</sub> desorption peak was detected before electron bombardment (dashed line). Electron bombardment of the NO<sub>2</sub> condensed surface for 10 s (7  $\mu$ A integrated current flux, 100 eV energy) leads to a dramatic decrease of the NO<sub>2</sub> multilayer peak and an apparent O<sub>2</sub> desorption peak evolved at 550 K (solid line), indicating the dissociation of NO<sub>2</sub> and deposition of oxygen. The heating rate used in these experiments was ~10 K/s.

#### **Results and Discussions**

There is no detectable thermal dissociation of NO<sub>2</sub> on Au-(111) based on TPD experiments (Figure 1). NO<sub>2</sub> is the only species evolved from the surface, based on a quantitative analysis of mass spectrometer fragmentation patterns. Specifically, there is no O<sub>2</sub> (32 amu) detected (Figure 1, dashed line). There are two NO<sub>2</sub> desorption peaks centered at 136 and 220 K, which are assigned to sublimation of NO<sub>2</sub> multilayers and desorption of NO<sub>2</sub> adsorbed directly on the surface. The multilayer peak grows indefinitely with exposure. These desorption temperatures are close to the previous TPD study of NO<sub>2</sub> on Au(111), where multilayer and monolayer NO<sub>2</sub> were found to desorb at 143 and 230 K, respectively.<sup>12,17</sup> The lack of NO<sub>2</sub> dissociation on clean Au(111) is also consistent with previous studies.<sup>12</sup>

NO<sub>2</sub> exposure (1  $\times$  10<sup>-10</sup> Torr pressure rise for 1 min) followed by electron bombardment (integrated current flux of 7  $\mu$ A with 100 eV energy) leads to a decrease of the NO<sub>2</sub> (46 amu) peak and an increase of the  $O_2$  (32 amu) peak, as shown in Figure 1 (solid line). Specifically, 10 s of electron bombardment results in a dramatic decrease of multilayer NO<sub>2</sub>, while monolayer NO<sub>2</sub> remains intact. Simultaneously, an O<sub>2</sub> desorption peak was detected at 550 K. Comparing with other studies,<sup>14</sup> this O<sub>2</sub> desorption feature is readily assigned to the recombination of atomic oxygen on the surface, indicating that atomic oxygen was deposited via the electron-induced dissociation of NO<sub>2</sub>. No evidence for deposition of nitrogen was found in the experiments, based on the absence of N<sub>2</sub> production in our desorption data. The only other species detected in our temperature programmed desorption experiments were NO and NO<sub>2</sub>. We also found that deposition of atomic oxygen through electron bombardment competes with electron induced desorption of NO<sub>2</sub>, yielding the depletion of NO<sub>2</sub> and thereby limiting the coverage of atomic oxygen.

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*Figure 2.* (a) Temperature programmed desorption of O<sub>2</sub> from Au(111) through the recombination of atomic oxygen for various initial coverages of O and (b) the corresponding absolute atomic oxygen coverage as a function of the number of cycles of NO<sub>2</sub> exposure and electron bombardment (note: the linear lines are fits by the program). The absolute coverage is estimated using O(1s) X-ray photoelectron spectra (data not shown). The heating rate used in all TPD experiments was ~10 K/s.

Iterative NO<sub>2</sub> exposing and electron bombardment were carried out to prepare higher coverages of atomic oxygen on Au(111) surface until saturation coverage was reached. Briefly, NO<sub>2</sub> was directly exposed at a pressure rise of  $1 \times 10^{-10}$  Torr at 100 K, followed by the electron bombardment (integrated current flux of 7  $\mu$ A with 100 eV energy) for 1 min, leading to atomic oxygen deposition from NO<sub>2</sub>. Repetition of these processes resulted in an increase in the oxygen coverage. Figure 2a shows O<sub>2</sub> TPD data with respect to the number of repetitive cycles mentioned above. The intensity of the O<sub>2</sub> desorption peak increases as the number of cycles is increased and becomes constant after seven cycles.

The absolute coverage of atomic oxygen deposited on Au-(111) was estimated by comparing the XPS intensity ratio (O(1s)/Au(4f)) of saturated oxygen atoms with that measured for chemisorbed NO<sub>2</sub> (data not shown). The coverage of the first layer NO<sub>2</sub> on Au(111) surface has been previously estimated by Koel and co-workers to be 0.4 monolayer.<sup>14</sup> Based on their calibration, the maximum coverage of oxygen atoms deposited using our method was calculated to be ~0.4 monolayer. This value is close to the saturation coverage of atomic oxygen deposited from the reaction of NO<sub>2</sub> and H<sub>2</sub>O ( $\theta_{\rm O} = 0.42$  monolayer),<sup>17</sup> but less than that obtained from ozone exposure ( $\theta_{\rm O} = 1.2$  monolayer).<sup>14</sup> The absolute coverage for different dosing and bombarding cycles has been determined from the ratio of the integrated area under the O<sub>2</sub> TPD peaks to that of the saturation coverage. Figure 2b shows that the increase of atomic oxygen coverage is nearly linear with repetitive cycles up to 7, at which point there is no further increase.

The peak temperature for O<sub>2</sub> evolution is 550 K, independent of oxygen coverage, implying pseudo-first-order kinetics of oxygen desorption. The apparent activation energy of desorption calculated assuming a preexponential factor of  $1 \times 10^{13} \text{ s}^{-1}$  and first-order kinetics<sup>22</sup> is ~33 kcal/mol, consistent with previous estimates of 31.7–33.6 kcal/mol.<sup>9,14</sup>

We attribute the pseudo-first-order kinetics of O<sub>2</sub> desorption peaks on Au(111) to a restructuring of the surface, based on LEED and STM data (described elsewhere<sup>18</sup>). Since the surface structure changes when oxygen is deposited from electron bombardment of NO<sub>2</sub>, we propose that the activation energy and/or preexponential factor for desorption of O2 are dependent on surface structure. This conclusion also raised the possibility that specific surface structures are required for O<sub>2</sub> dissociation and that the kinetics for  $O_2$  dissociation may be dependent on the method used for oxygen deposition. Alternative explanations for the observed pseudo-first-order kinetics are that the ratelimiting step for the  $O_2$  desoprtion on Au(111) is the conversion of an oxidic oxygen species to chemisorbed oxygen<sup>14</sup> or that attractive interactions control the rate, as proposed for Ni(110) where pseudo-first-order kinetics are observed when H<sub>2</sub> forms from one-dimensional chains of H atoms.<sup>24</sup> Our explanation is most similar to the latter case of attractive interactions because both resemble a surface phase transition.

Since atomic oxygen recombination and desorption is the microscopic reverse of  $O_2$  dissociation and adsorption, the surface restructuring observed after O deposition from electron bombardment of NO<sub>2</sub> may indicate a coverage-dependent enthalpy ( $\Delta E_{dis}$ ) and/or entropy ( $\nu_d$ ) change. Likewise, there may be a dependence in the kinetic barriers for O<sub>2</sub> dissociation and adsorption on oxygen coverage which lead to the dependence of O<sub>2</sub> dissociation probability on oxygen atom coverage. Thus, we studied the effect of oxygen pre-coverage on the dissociation probability of O<sub>2</sub> by performing studies of <sup>18</sup>O<sub>2</sub> at a surface temperature of 400 K.

As anticipated, there is no detectable dissociation of  ${}^{18}O_2$  on the clean Au(111)-herringbone surface after exposure to a background pressure of  $1 \times 10^{-7}$  Torr of  ${}^{18}O_2$  for 1 min, based on the absence of the signal for any isotope of dioxygen in TPD experiments (Figure 3, dashed lines). Specifically, there was no detectable 34 ( ${}^{16}O{}^{18}O$ ) and 36 ( ${}^{18}O_2$ ) evolution in TPD data obtained after exposing clean Au(111) to  ${}^{18}O_2$ .

Dissociation of  ${}^{18}O_2$  is detected after exposure of Au(111) precovered with atomic oxygen. Specifically, there is a clearly defined mass 34 ( ${}^{16}O^{18}O$ ) peak at 550 K detected after exposing  ${}^{18}O_2$  to Au(111) containing some  ${}^{16}O$  formed from NO<sub>2</sub> dissociation (Figure 3, solid lines). A peak associated with formation of  ${}^{16}O_2$  (32 amu) is observed coincident with the

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*Figure 3.* Temperature programmed desorption of <sup>16</sup>O<sub>2</sub>, <sup>16</sup>O<sup>18</sup>O, and <sup>18</sup>O<sub>2</sub> from Au(111). A small amount of atomic oxygen (0.1 monolayer) is deposited on Au(111) by 1 cycle of NO<sub>2</sub> dosing and electron bombardment before dosing <sup>18</sup>O<sub>2</sub> at 400 K. Note that <sup>18</sup>O<sub>2</sub> was dosed using background dosing with a pressure of  $1 \times 10^{-7}$  Torr for 1 min. The appearance of desorption peak at 550 K with m/z = 34 indicates the dissociation of <sup>18</sup>O<sub>2</sub> on atomic oxygen covered Au(111) (solid lines). <sup>18</sup>O<sub>2</sub> was also dosed on clean Au(111) under the same conditions, in which case no apparent dissociation was observed (dashed lines).



*Figure 4.* Dissociation probability of <sup>18</sup>O<sub>2</sub> on a Au(111) as a function of oxygen precoverage. The oxygen precovered Au(111) was prepared via electron bombardment of condensed NO<sub>2</sub>, and <sup>18</sup>O<sub>2</sub> was exposed at 400 K using background dosing with a pressure of  $1 \times 10^{-7}$  Torr for 1 min for all experiments. The dissociation probability is estimated from the temperature programmed desorption experiments, as described in the experimental section.

<sup>16</sup>O<sup>18</sup>O. No 36 (<sup>18</sup>O<sub>2</sub>) is detected, which is consistent with the large amount of <sup>16</sup>O present on the surface, based on the ratio of the 32 amu/34 amu signal that indicates an abundance of <sup>16</sup>O relative to <sup>18</sup>O, yielding an estimate of <sup>18</sup>O:<sup>16</sup>O < 1:20. Analogous experiments were also performed in a separate vacuum chamber with a base pressure of  $< 3 \times 10^{-10}$  Torr and using a separate Au(111) crystal. The qualitatively reproduced <sup>18</sup>O<sub>2</sub> dissociation results rule out artifacts in the data shown.

The <sup>18</sup>O<sub>2</sub> dissociation probability on Au(111) as a function of oxygen precoverage was estimated using quantitative analysis of the TPD experiments (Figure 4). Specifically, <sup>18</sup>O<sub>2</sub> dissociation probability on clean Au(111) is less than  $10^{-6}$  (estimated from the noise level of the mass spectrometer). When predepositing small amount of atomic oxygen (0.02 monolayer), the <sup>18</sup>O<sub>2</sub> dissociation probability significantly increases to  $10^{-3}$ . The <sup>18</sup>O<sub>2</sub> dissociation probability increases at low oxygen precoverages (up to 0.1 monolayer). At higher oxygen precoverages, the dissociation probability becomes nearly constant, reaching a value of  $2 \times 10^{-3}$  for oxygen precoverages in the range of 0.1–0.4 monolayer.

The significant enhancement of O2 dissociation on the oxygen precovered Au(111) surface may be due to the release of Au atoms from the herringbone reconstructed Au(111). In a companion paper, we show that adsorption of atomic oxygen using electron-induced dissociation of NO<sub>2</sub> induces restructuring of the surface and release of Au atoms from the surface.<sup>18</sup> Indeed, theoretical studies by Mills et al. showed small Au clusters on Au(111) have a higher binding energy with  $O_2$ compared to a Au(111) surface,<sup>19</sup> as a result of having a large number of low-coordinated Au atoms which facilitate orbital overlapping between Au and O<sub>2</sub>. Our recent STM study of oxygen deposition on Au(111) shows restructuring of Au(111)herringbone structure accompanied by the formation of small gold island and/or serrated step edge due to the release of gold atoms at the expense of herringbone elbow sites (dislocation sites).<sup>18</sup> These results support the possibility that the creation of low-coordinated gold atoms is responsible for the enhancement of O<sub>2</sub> dissociation.

Accordingly, the high activity of small Au clusters on metaloxide supports for CO oxidation might originate from the enhanced  $O_2$  dissociation due to the large number of lowcoordinated Au atoms. Considering that  $O_2$  dissociation is most likely the key, rate-limiting step in catalytic oxidation processes (as supported by the fact that Au(111) or Au(110) can also serve as oxidation catalysts once oxygen is deposited on these surfaces<sup>7,9</sup>), it is not surprising that small Au clusters, which can enhance  $O_2$  dissociation, show high catalytic activity in oxidation. Introducing Au atoms on a Au surface via physical vapor deposition (PVD) may provide more compelling evidence for all the discussions above, and these studies are currently under investigation in our lab.

#### Conclusions

The dissociation probability of O<sub>2</sub> on Au(111) at 400 K significantly increases up to  $10^{-3}$  when a small amount of atomic oxygen is present on the surface. The enhancement of O<sub>2</sub> dissociation is possibly due to the release of Au atoms from the surface restructuring induced by atomic oxygen deposition. Atomic oxygen is deposited by electron induced NO2 dissociation on the surface at 100 K. A saturation coverage is attained by the repetition of NO<sub>2</sub> dosing and electron bombardment. Desorption of O<sub>2</sub>, coming from recombination of atomic oxygen on Au(111), occurs at 550 K and does not shift with coverage, indicating pseudo-first-order kinetics. We attribute this unusual kinetics to the surface restructuring by oxygen deposition, which leads to a coverage-dependent enthalpy ( $\Delta E_{dis}$ ) and/or entropy  $(\nu_d)$ . In addition, the surface restructuring and the coveragedependent enthalpy ( $\Delta E_{dis}$ ) and/or entropy ( $\nu_d$ ) may also serve as a kinetic barrier for  $O_2$  dissociation on Au(111), resulting in an extremely low dissociation probability of O2 on clean Au-(111) (less than  $10^{-6}$ ).

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